

Key concept in Stat. Mech. → Thermodynamic properties can exhibit fluctuations around their mean value.

- what is their magnitude?
- consequences on behavior?

$$P_m = \frac{e^{-\beta E_m}}{Q(T, V, N)} \quad \text{in canonical } (N, V, T) \text{ ensemble}$$

$$\langle X \rangle = \sum_{\substack{\text{all states} \\ \text{at } N, V}} X_m P_m \quad \leftarrow \text{average value of } X$$

$$\begin{aligned} \sigma_x^2 &= \langle (X - \langle X \rangle)^2 \rangle = \langle X^2 - 2X\langle X \rangle + \langle X \rangle^2 \rangle \\ &= \langle X^2 \rangle - 2\langle X \rangle \langle X \rangle + \langle X \rangle^2 = \langle X^2 \rangle - \langle X \rangle^2 \end{aligned}$$

$$\Rightarrow \sigma_x^2 = \sum P_m X_m^2 - (\sum P_m X_m)^2$$

### Energy distribution and fluctuations

$$\begin{aligned} P(E) &= \sum_m P_m \delta_{E_m, E} = \sum_m \frac{e^{-\beta E_m}}{Q(T, V, N)} \delta_{E_m, E} = \\ &= \frac{Q(E, V, N) e^{-\beta E}}{Q(T, V, N)} \end{aligned}$$

$Q$  can also be written as a summation over energies (rather than microstates):

$$Q = \sum_{\text{all } E_i} Q(E_i, V, N) e^{-\beta E_i} \quad E_i: \text{energy levels}$$

Operation on partition function is a Laplace Transform (microscopic analog of Legendre transform)

Fluctuations:  $\sigma_E^2 = \langle E^2 \rangle - \langle E \rangle^2 = \sum_{\text{microstates}} P_m E_m^2 - \left( \sum_m P_m E_m \right)^2$

$$Q = \sum_m e^{-\beta E_m}; P_m = e^{-\beta E_m} / Q$$

$$\frac{\partial Q}{\partial \beta}_{V,N} = - \sum_m E_m e^{-\beta E_m} = - Q \sum_m P_m E_m$$

$$\frac{\partial^2 Q}{\partial \beta^2}_{V,N} = \sum_m E_m^2 e^{-\beta E_m} = Q \sum_m P_m E_m^2$$

$$\sigma_E^2 = \frac{1}{Q} \frac{\partial^2 Q}{\partial \beta^2} - \left[ \frac{1}{Q} \frac{\partial Q}{\partial \beta} \right]^2 = \frac{\partial^2 \ln Q}{\partial \beta^2} = - \frac{\partial \langle E \rangle}{\partial \beta}_{V,N}$$

$\therefore$  Fluctuations are proportional to the second derivative of the ln [relevant partition function] (Thermodynamic potential)

$$\ln \Omega = \frac{S}{k_B} = y^{(0)} \quad d\ln \Omega = \beta dE + \beta P dV - \beta \mu dN$$

$$\ln Q = -\frac{A}{k_B T} = y^{(1)} \quad d\ln Q = -E dB + \beta P dV - \beta \mu dN$$

$$\sigma_E^2 = \langle (\delta E)^2 \rangle = \langle (E - \langle E \rangle)^2 \rangle = \frac{\partial^2 \ln Q}{\partial \beta^2} = - \frac{\partial E}{\partial \beta} = \underline{k_B T^2 C_V}$$

The mean-square fluctuation of  $E$  is proportional to the heat capacity at const. volume.

Relative magnitude  $\frac{\sqrt{\sigma_E^2}}{E} = \frac{\sqrt{C_V k_B T^2}}{E} \propto \frac{\sqrt{N}}{N} = \frac{1}{\sqrt{N}}$

Relative fluctuations go down as  $\sqrt{N}$

$$N \approx 1000 \text{ (simulations)} \quad \frac{\sqrt{\sigma_E^2}}{E} \sim 3\% \\ N \approx 10^{23} \text{ (lab)} \quad \frac{\sqrt{\sigma_E^2}}{E} \sim 10^{-11}$$

The relationship  $\sigma_E^2 = \frac{\partial^2 \ln Q}{\partial \beta^2} \Big|_{V,N}$

for  $\langle E^2 \rangle - \langle E \rangle^2$  in the NVT ensemble can be readily generalized to any ensemble -

E.g. grand Canonical ( $\mu \text{VT}$ ) ensemble:

$$\langle N^2 \rangle - \langle N \rangle^2 = \sigma_N^2 = \left( \frac{\partial^2 \ln \Xi}{\partial (\beta \mu)^2} \right)_{B,V} = k_B T \left[ \frac{\partial \langle N \rangle}{\partial \mu} \right]_{T,V}$$

Constant-pressure (NPT) ensemble:

$$\langle V^2 \rangle - \langle V \rangle^2 = \sigma_V^2 = \left( \frac{\partial^2 \ln \Delta}{\partial (\beta P)^2} \right)_{B,N} = - \left( \frac{\partial \langle V \rangle}{\partial (\beta P)} \right)_{B,N} = -k_B T \left[ \frac{\partial \langle V \rangle}{\partial P} \right]_{T,N}$$

Note that these same derivatives in classical thermodynamics correspond to stability conditions.

Near a limit of stability (or near a critical point), the corresponding fluctuations diverge.

These can even be generalized for cross-fluctuations, e.g. in  $\mu \text{VT}$  ensemble:

$$\begin{aligned} \langle \delta E \delta N \rangle &= \langle (E - \langle E \rangle)(N - \langle N \rangle) \rangle = \langle \delta E \rangle \langle \delta N \rangle \\ &= \frac{\partial^2 \ln \Xi}{\partial (\beta \mu) \partial \beta} = \left( \frac{\partial \langle N \rangle}{\partial \beta} \right)_{T,V} = -k_B T^2 \left( \frac{\partial \langle N \rangle}{\partial T} \right)_{\mu,V} \end{aligned}$$