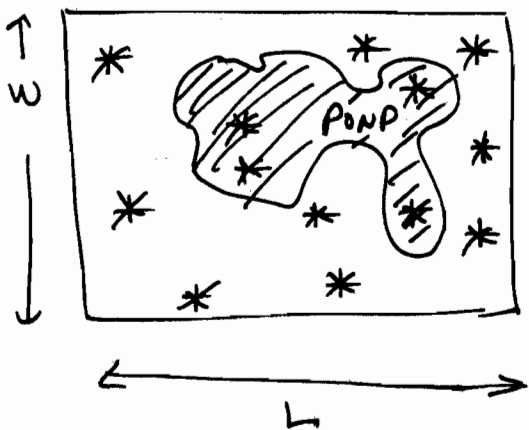


History

von Neumann, Ulam + Metropolis (1947) → studies of diffusion of neutrons in fissionable matter. Coined the term "Monte Carlo".

Metropolis, Rosenbluth², Teller² (1953) → first MC simulation for equilibrium stat. mech. properties.

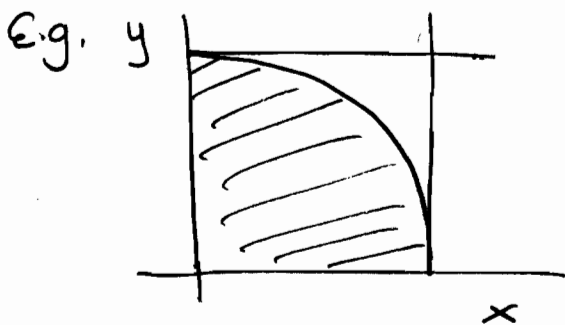
Monte Carlo Integration



Consider a pond of arbitrary shape. Possible way: throw rocks at random, in area bounded by $w \times L$ rectangle

$$\text{Area of pond} = w \times L \times \frac{\text{hits in pond}}{\text{total hits}}$$

Finding an area is an integration problem, so that "Monte Carlo" is a way to perform multidimensional integrals.



Find area of quarter circle (or value of π) by above method:

After 10^7 trials $\pi \approx 3.1417$
(4 sig. figures)

The gain in accuracy is only 1 decimal point per $10\times$ in the number of trials.

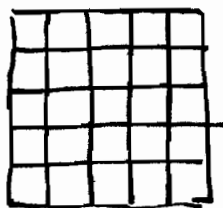
For a system of low dimensionality, analytical or direct numerical methods (e.g. Simpson's rule) work much better than Monte Carlo. E.g. for 10^4 function evaluations (intervals) Simpson's rule integration gives $\pi \approx 3.141593$ (6 sig. figures)

Multidimensional Integrals

Canonical Partition function for $N \approx 20$ particles

$$Q = \sum_{\text{pos } 1} \sum_{\text{pos } 2} \sum_{\text{pos } 3} \dots \sum_{\text{pos } N} \exp(-\beta U)$$

Even if each particle only has 25 possible positions, one needs $25^{20} \approx 10^{28}$ function evaluations for a direct integration. Moreover, most of the configurations contribute very little to Q :



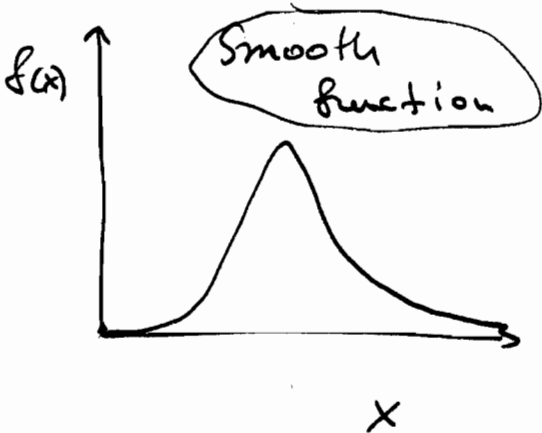
← 25 positions

① ② ③ ④ ⑤ ⑥ ⑦ ⑧ ⑨ ⑩ ... ⑳

Throw 20 particles at random

Assume no two particles can occupy same position in space. What fraction of configurations generated this way are allowable?

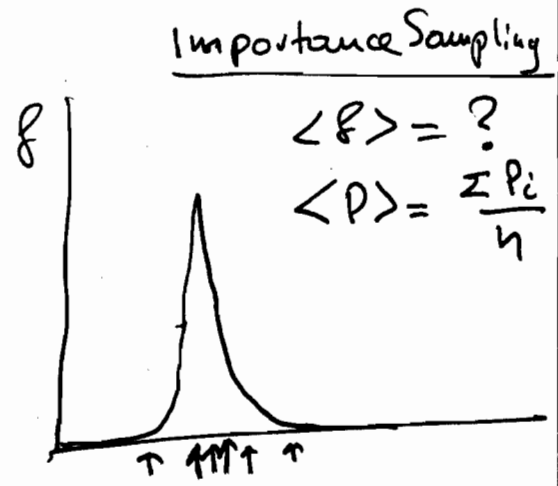
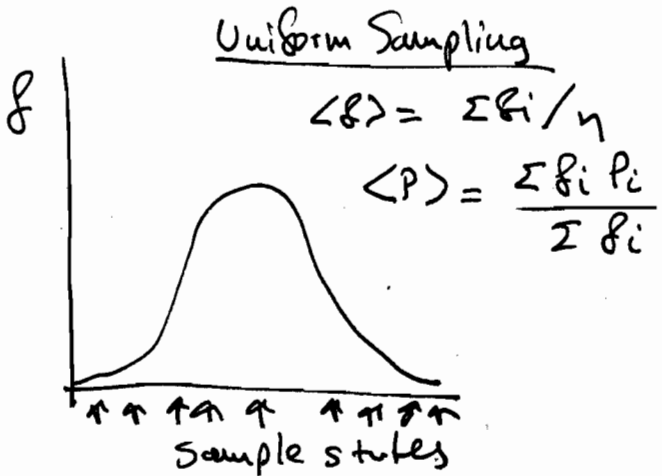
Things are better for less dense systems, but most of the systems of interest in stat. mechanics have this problem \rightarrow function to be integrated has non-zero values only in small regions of phase space



Importance Sampling (also called Metropolis's MC)

Overcomes sampling problem by generating states with a desired probability distribution, rather than uniformly.

This does not work for getting $\int f(x) dx$, but does work for getting properties that depend on $f(x)$:



How do we generate states with the correct probability?

Consider the system of 5 particles in 10 positions - the number of possible states (for distinguishable particles) is $10^5 = 100,000$ states

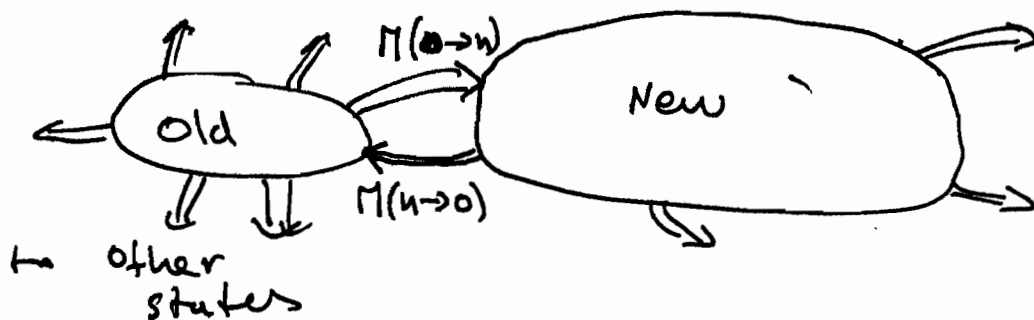
Now consider two configurations Old (o)
and New (n)

In a thought experiment, we can perform a MC simulation with a number of samples much greater than the number of states (samples $\gg 10^5$). In this simulation, we generate uniform random configurations. In the large sample, there are $N(o)$ and $N(n)$ realizations for the "old" and "new" configurations, with

$$\frac{N(o)}{N(n)} = \frac{\exp(-\beta U(o))}{\exp(-\beta U(n))} = \exp(-\beta \Delta U) \quad \left[\begin{array}{l} \text{in the} \\ \text{NVT} \\ \text{ensemble!} \end{array} \right]$$

Now consider the case of non-uniform sampling.

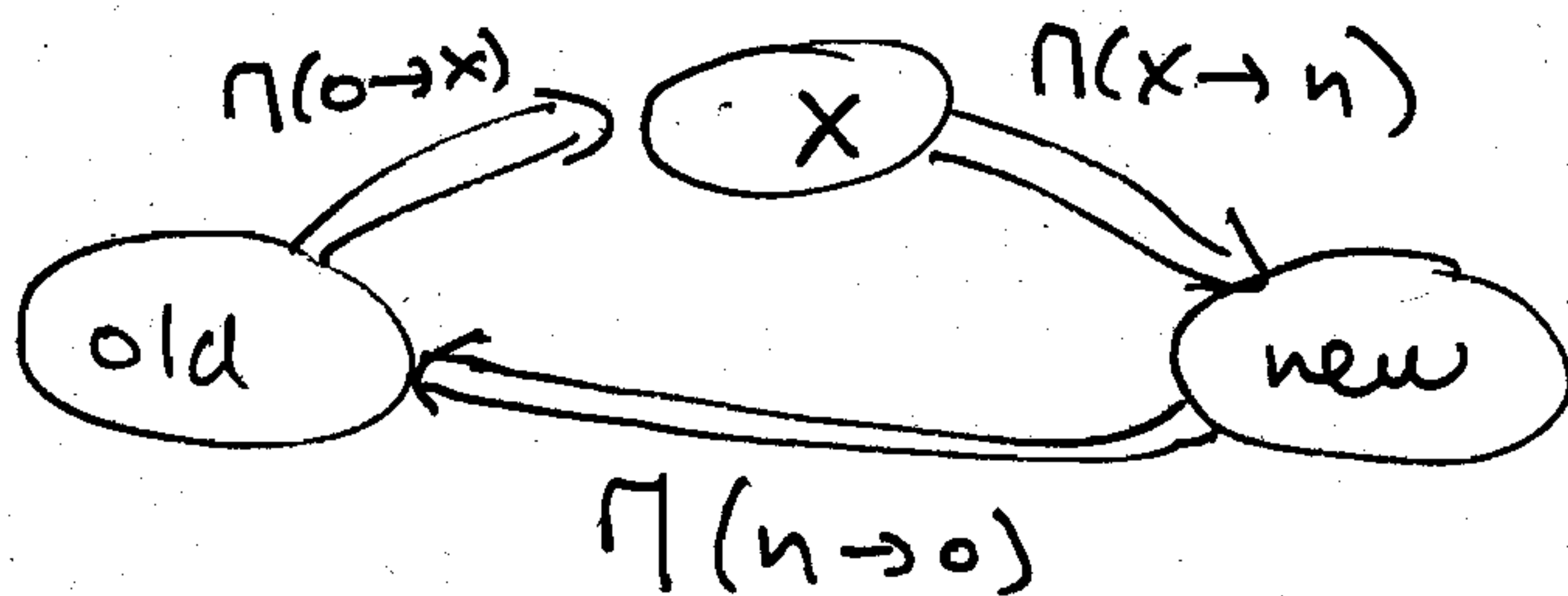
To ensure that configurations appear with the correct probability, one possible way is to require that in the evolution of configurations,



Detailed Balance Condition: $N(o) \cdot \Pi(o \rightarrow n) = N(n) \cdot \Pi(n \rightarrow o)$

In other words, we must have that the probabilities of going from $o \rightarrow n$ and $n \rightarrow o$ must be inversely proportional to the desired probability of finding each state.

Why is this only one possible way? \rightarrow One could have this:



See Manousiouthakis & Deem, J. Chem. Phys. 110:2753 (1999) for proof Detailed Balance is not necessary

But, in practice, it is very hard to make sure that this type of simulation (that violates detailed balance) is correct.

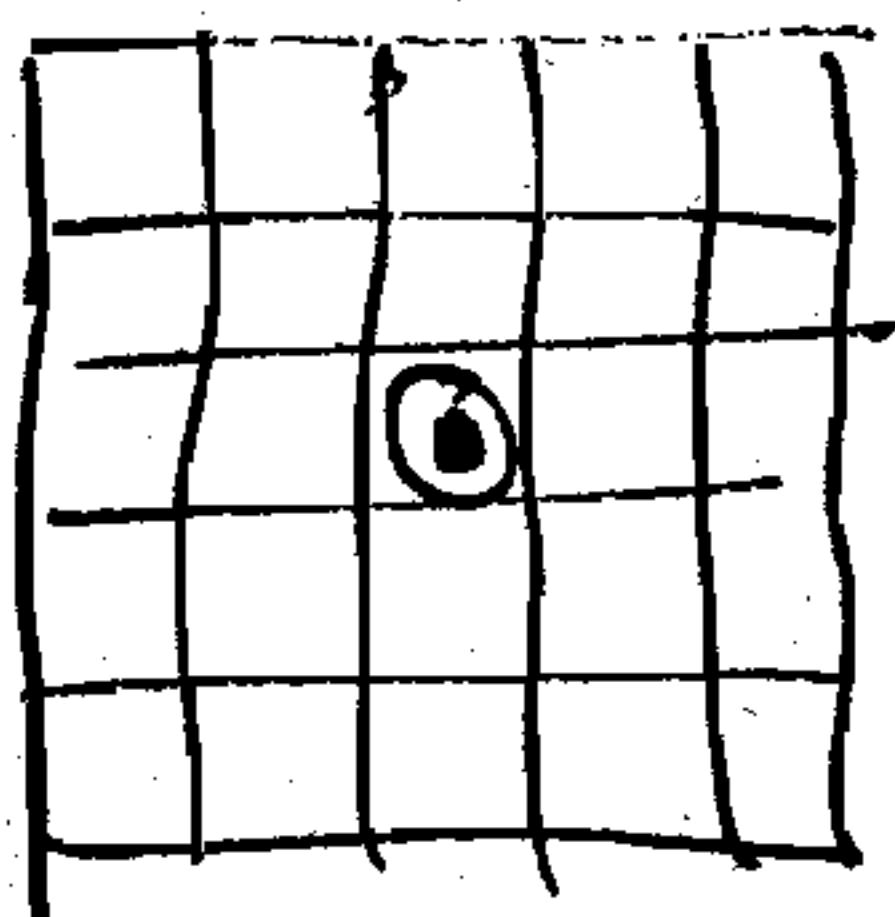
The probability of going from state o to state n has two parts:

$$\Pi(o \rightarrow n) = \alpha(o \rightarrow n) \times \text{accept}(o \rightarrow n)$$

\uparrow probability of selecting State n from state o
 \uparrow to be determined

$\alpha(o \rightarrow n)$ contains the rules of the simulation -

e.g.



{ try to move particle to a new position to the left, right, top, or bottom of the old position }

The best choice for $\alpha(0 \rightarrow n)$ is one that satisfies

microscopic reversibility $\alpha(0 \rightarrow n) = \alpha(n \rightarrow 0)$

so that the probability of trying state n from 0 is the same as the probability of trying 0 from n .

E.g. attempted move ① select random particle, attempt to move it to the left or top position
 \hookrightarrow updates microscopic reversibility

- ② select random particle, attempt to move to a randomly selected position \rightarrow OK
- ③ select random particle, attempt to move it to a randomly selected unoccupied position \rightarrow is this OK?
 [class problem]

Now we can derive the acceptance condition:

$$N(0) \cdot \alpha(0 \rightarrow n) = N(n) \cdot \alpha(n \rightarrow 0) \Rightarrow$$

$$N(0) \cdot \alpha(0 \rightarrow n) \cdot \text{accept}(0 \rightarrow n) = N(n) \cdot \alpha(n \rightarrow 0) \cdot \text{accept}(n \rightarrow 0)$$

$$\Rightarrow \frac{\text{acc}(0 \rightarrow n)}{\text{acc}(n \rightarrow 0)} = \frac{N(n)}{N(0)} = \exp(-\beta [u(n) - u(0)])$$

There are many possibilities for $\text{acc}(0 \rightarrow n)$ that satisfy this equation - for example:

Metropolis's Rule:

$$\begin{aligned} \text{acc}(0 \rightarrow n) &= 1 && \text{if } u(n) < u(0) \\ \text{acc}(0 \rightarrow n) &= \exp(-\beta[u(n) - u(0)]) && \text{if } u(n) \geq u(0) \end{aligned}$$

Symmetrical (Barker)

Rule:

$$\text{acc}(0 \rightarrow n) = \frac{\exp(-\beta u(n))}{\exp(-\beta u(0)) + \exp(-\beta u(n))}$$

Example of application of these rules:

Consider a system with exactly 10 states,

numbered j :

1	2	3	4	5	6	7	8	9	10
1	2	3	4	5	6	7	8	9	10

 } probability is equal to index

What is the partition function? (Class Problem)
 What is the equilibrium probability of state 1?

Now, let's make an importance sampling simulation.

Transition rule? $\alpha(0 \rightarrow n) = 1/10$ for any pair of states 0 and n (including $0=n$)

$\Pi(0 \rightarrow n)$ Metropolis's

		$\rightarrow n$								
\downarrow		1	2	3	...	10				
0	1	1/10	1/10	1/10	...	1/10				
	2	1/20	3/20	1/10	...	1/10				
	3	1/30	2/30	1/10				
	...									
	10	1/100	2/100	11/20				

$\Pi(0 \rightarrow n)$ Barker

		$\rightarrow n$								
\downarrow		1	2	3	...	10				
0	1	0.252	2/30	3/40	...	10/110				
	2	1/30	0.371	3/50	...	10/120				
	3	1/40								
	...									
	10	1/110				0.719				



Barker's probability of remaining in the same state is higher than for Metropolis's acceptance. In general, convergence from an arbitrary starting state to equilibrium is faster with Metropolis's acceptance.

Example of flow of probabilities with the system with 10 states - The starting state is state 3, which we represent as $[0\ 0\ 1\ 0\ 0\ 0\ 0\ 0\ 0\ 0]$

Trial

Metropolis's

0	0	0	1	0	0	0	0	0	0	0
1	0.033	0.067	0.200	0.100	0.100	0.100	0.100	0.100	0.100	0.100
2	0.024	0.049	0.083	0.089	0.102	0.114	0.124	0.132	0.139	0.145
3	0.02	0.04	0.062	0.079	0.097	0.113	0.128	0.142	0.154	0.165
4	0.019	0.038	0.057	0.075	0.093	0.111	0.128	0.145	0.16	0.174
5	0.018	0.037	0.055	0.074	0.092	0.11	0.128	0.145	0.162	0.178
6	0.018	0.037	0.055	0.073	0.091	0.109	0.128	0.145	0.163	0.18
7	0.018	0.036	0.055	0.073	0.091	0.109	0.127	0.146	0.163	0.181
8	0.018	0.036	0.055	0.073	0.091	0.109	0.127	0.145	0.164	0.182
∞	0.018	0.036	0.055	0.073	0.091	0.109	0.127	0.145	0.164	0.182

Barker

0	0	0	1	0	0	0	0	0	0	0
1	0.025	0.040	0.454	0.057	0.063	0.067	0.070	0.073	0.075	0.077
2	0.025	0.046	0.225	0.075	0.086	0.095	0.103	0.109	0.115	0.12
3	0.023	0.044	0.128	0.079	0.093	0.106	0.118	0.128	0.137	0.145
4	0.021	0.041	0.087	0.078	0.095	0.11	0.124	0.137	0.149	0.16
5	0.019	0.039	0.069	0.077	0.094	0.111	0.127	0.141	0.155	0.168
6	0.019	0.038	0.061	0.075	0.093	0.111	0.128	0.144	0.159	0.173
7	0.019	0.037	0.057	0.074	0.093	0.11	0.128	0.145	0.161	0.176
8	0.018	0.037	0.056	0.074	0.092	0.11	0.128	0.145	0.162	0.178
9	0.018	0.037	0.055	0.073	0.092	0.11	0.128	0.145	0.163	0.18
10	0.018	0.036	0.055	0.073	0.091	0.11	0.128	0.145	0.163	0.18
11	0.018	0.036	0.055	0.073	0.091	0.109	0.127	0.145	0.163	0.181
12	0.018	0.036	0.055	0.073	0.091	0.109	0.127	0.146	0.163	0.181
13	0.018	0.036	0.055	0.073	0.091	0.109	0.127	0.145	0.164	0.181
14	0.018	0.036	0.055	0.073	0.091	0.109	0.127	0.145	0.164	0.182
∞	0.018	0.036	0.055	0.073	0.091	0.109	0.127	0.145	0.164	0.182

22-141 50 SHEETS
22-142 100 SHEETS
22-144 200 SHEETS

